



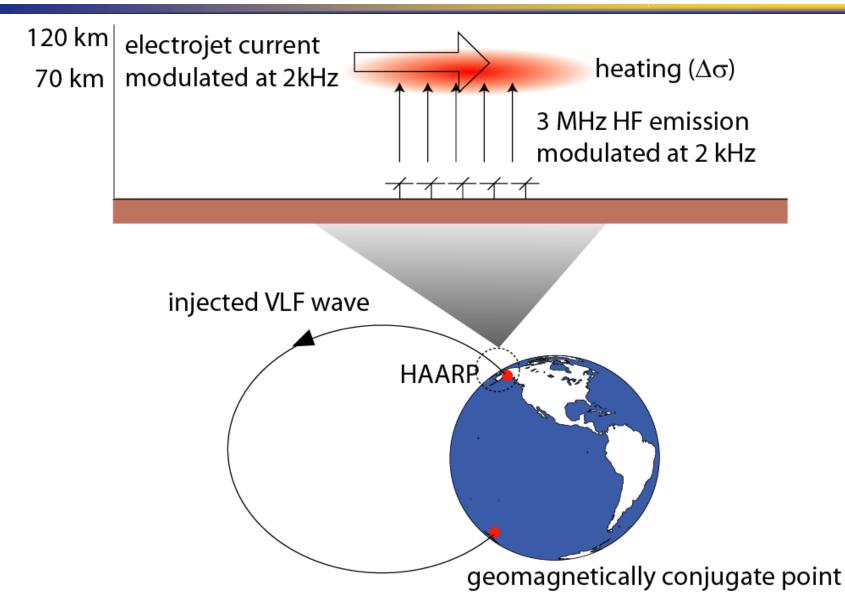
Ionospheric modification and ELF/VLF wave generation by HAARP

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Generation of ELF/VLF waves











High Frequency Active Auroral Research Program



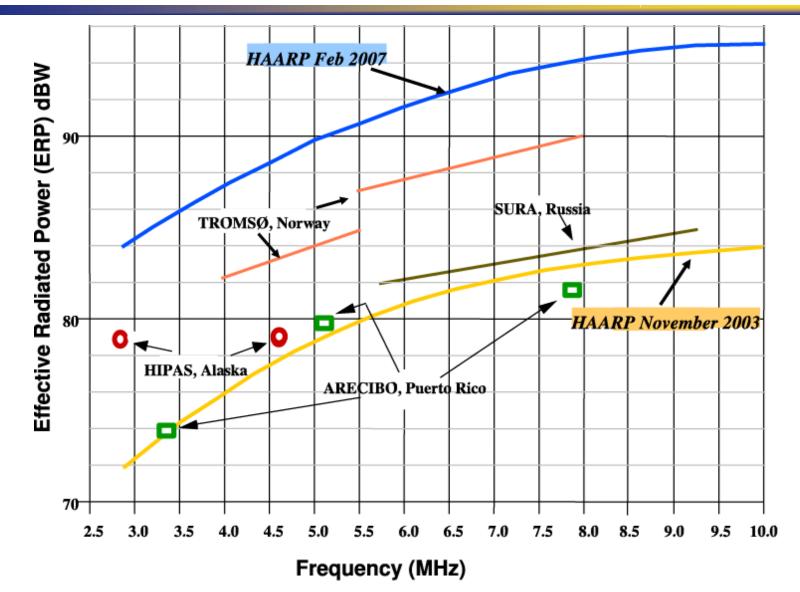
After upgrade in March 2006:

- 180 crossed dipole antennas
- 3.6 MW power
- ~2 GW effective radiated HF power (2.8-10 MHz) (lightning has ~20 GW isotropic ERP)



HAARP and other HF heating facilities

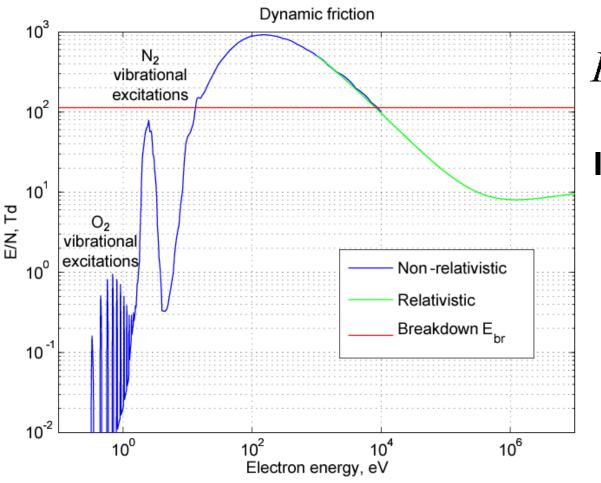






Important electron-molecule interaction concept: Dynamic friction force





$F = \sum N \sigma_i(v) \Delta \varepsilon_i$

Inelastic processes:

- Rotational, vibrational, electronic level excitations
- Dissociative losses
- Ionization

$(E/N)_{br}$ =130 Td where 1 Td = 10⁻²¹ V-m²



Kinetic Equation Solver (modified ELENDIF)



Time-dependent solution for $f(v,t) = f_0(v,t) + \cos\theta f_1(v,t)$ (almost isotropic)

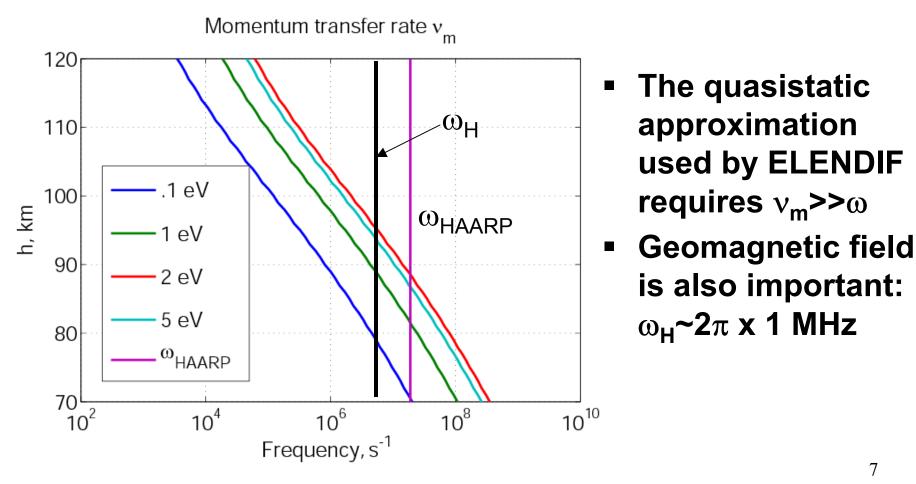
Physical processes inluded in ELENDIF:

- Quasistatic electric field
- Elastic scattering on neutrals and ions
- Inelastic and superelastic scattering
- Electron-electron collisions
- Attachment and ionization
- Photon-electron processes
- External source of electrons

New:

- Non-static (harmonic) electric field
- Geomagnetic field









Margenau distribution

$$f_0 = C \exp\left[-\frac{3m^3}{4Me^2E^2l^2}(v^4 + 2v^2\omega_{eff}^2l^2)\right]$$

where $l = v/v_m = (N\sigma_m)^{-1} = const$

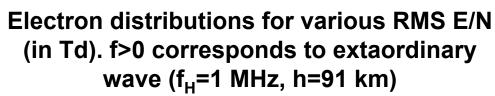
Druyvesteyn distribution ω=0

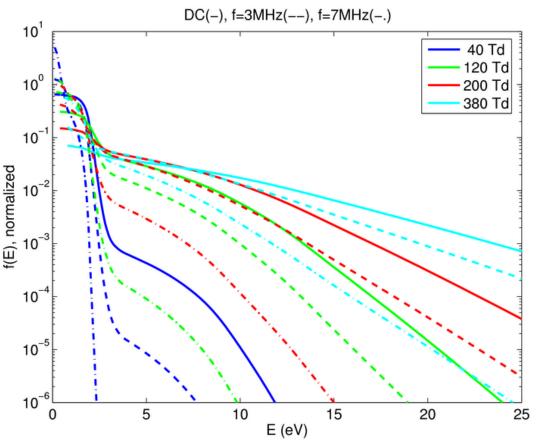
$$f_0 = C \exp\left[-\frac{3m^3v^4}{4Me^2E^2l^2}\right]$$



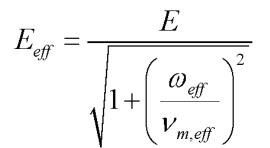
Calculated electron distributions







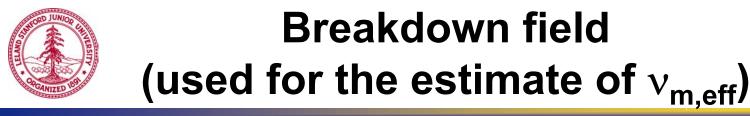
Effective electric field is smaller than in DC case:



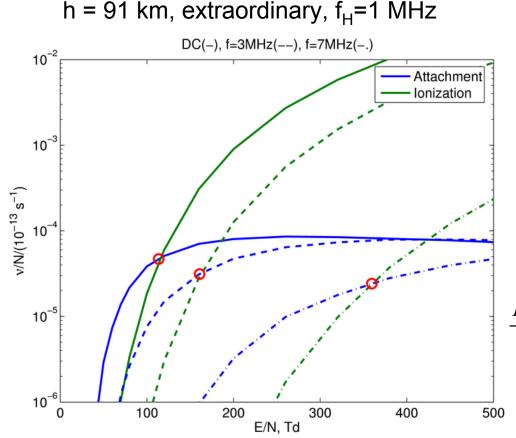
$$v_{m,eff} / N = 2 \times 10^{-13} \, s^{-1} m^3$$

 $\omega_{eff} = \omega \pm \omega_{H}$

+ ordinary - extraordinary







- Breakdown occurs when v_{ion}>v_{att}
- The point of breakdown (shown with •) shifts up in oscillating field

$$\frac{E_{br}}{N} \approx \left(\frac{E_{br}}{N}\right)_{DC} \sqrt{1 + \left(\frac{\omega \pm \omega_{H}}{N \times 2 \times 10^{-13} \, s^{-1} m^{3}}\right)^{2}}$$

 f(v) at ionization energy (~15 eV) is most important





Power flux (1D), including losses:

$$\frac{dS}{dz} = -\alpha(S, z)S$$
$$\alpha = 2 \operatorname{Im} k + \frac{2}{R}$$
$$k = \frac{\omega}{c} \sqrt{1 + \frac{i\sigma}{\omega\varepsilon_0}}$$

HF conductivity (ordinary/extaordinary)

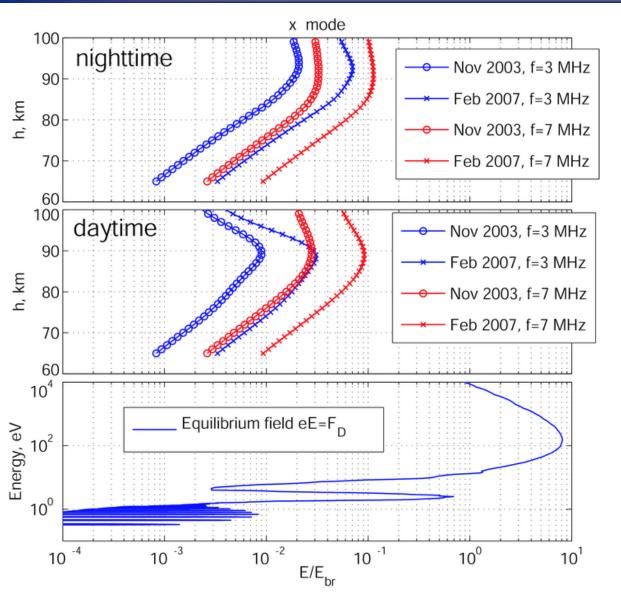
$$\sigma_{o,x} = -\frac{2e^2}{3m} \int \frac{\varepsilon^{3/2}}{V_m - i(\omega \pm \omega_H)} \frac{\partial}{\partial \varepsilon} \frac{n}{\varepsilon^{1/2}} d\varepsilon$$

11



Calculated HF electric field





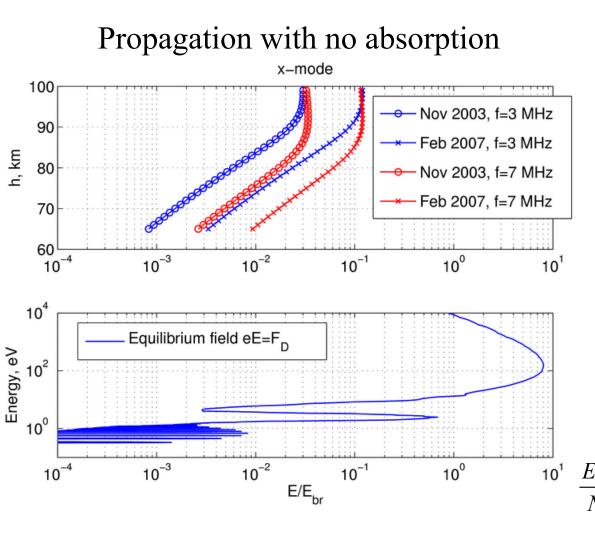
•Normalized field, E/E_{br} is shown

•For comparison, we show the dynamic friction function

•The N₂ vibrational threshold or breakdown field are not exceeded for current or upgraded HAARP power 12







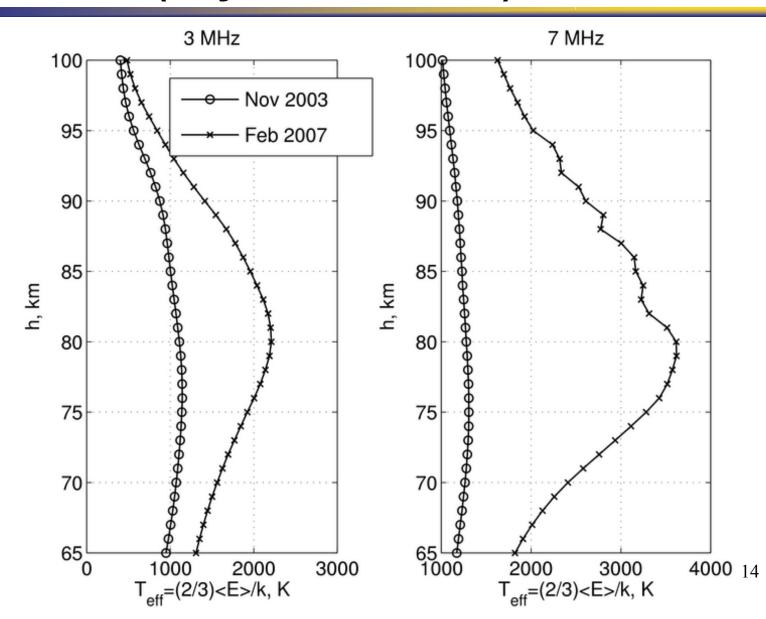
- The electric field can be higher in a non-steady state case
- Electric breakdown field with altitude:
 - Decreases due to thinning atmosphere
 - But, increases due to oscillations and magnetization.

 $\frac{10^{1}}{10^{1}} \frac{E_{br}}{N} \approx \left(\frac{E_{br}}{N}\right)_{DC} \sqrt{1 + \left(\frac{\omega \pm \omega_{H}}{N \times 2 \times 10^{-13} \, s^{-1} m^{3}}\right)}$



Temperature modification (daytime, x mode)

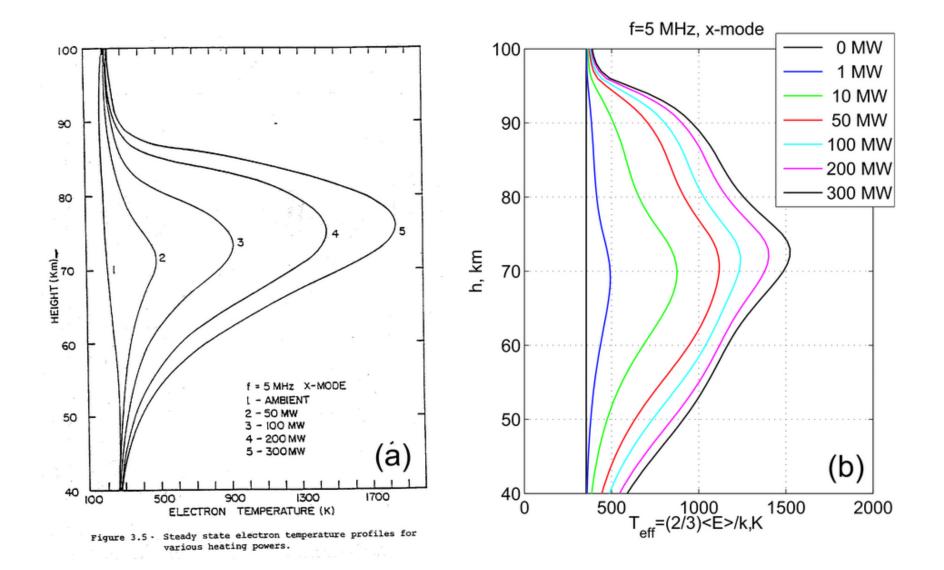






Comparison of Maxwellian and non-Maxwellian approaches









DC conductivity changes (for electrojet current)





- Conductivity changes due to modification of electron distribution
- Approximate formulas were used previously
- Pedersen (transverse)

$$\sigma_{p} = -\frac{2e^{2}}{3m} \int \frac{\nu_{m} \varepsilon^{3/2}}{\omega_{H}^{2} + \nu_{m}^{2}} \frac{\partial}{\partial \varepsilon} \frac{n}{\varepsilon^{1/2}} d\varepsilon \simeq \frac{N_{e} e^{2}}{m} \left\langle \frac{\nu_{m}}{\omega_{H}^{2} + \nu_{m}^{2}} \right\rangle$$

Hall (off-diagonal)

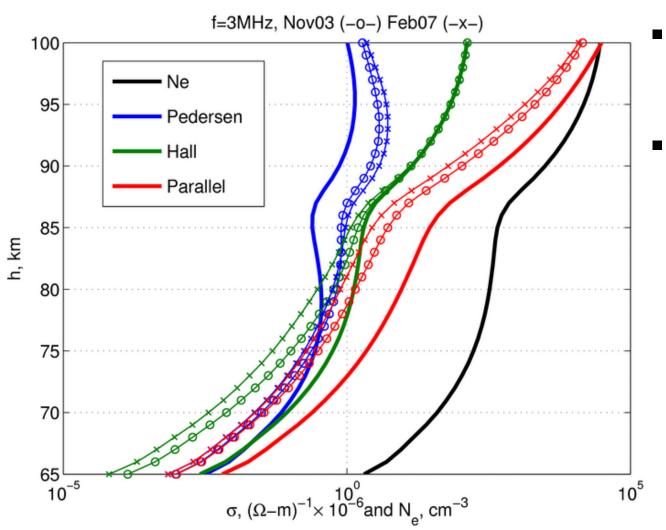
$$\sigma_{h} = -\frac{2e^{2}}{3m} \int \frac{\omega_{H} \varepsilon^{3/2}}{\omega_{H}^{2} + v_{m}^{2}} \frac{\partial}{\partial \varepsilon} \frac{n}{\varepsilon^{1/2}} d\varepsilon \simeq \frac{N_{e} e^{2}}{m} \left\langle \frac{\omega_{H}}{\omega_{H}^{2} + v_{m}^{2}} \right\rangle$$

• Parallel $\sigma_{z} = -\frac{2e^{2}}{3m} \int \frac{\varepsilon^{3/2}}{v_{m}} \frac{\partial}{\partial \varepsilon} \frac{n}{\varepsilon^{1/2}} d\varepsilon \simeq \frac{N_{e}e^{2}}{m} \left\langle \frac{1}{v_{m}} \right\rangle \qquad 17$



Conductivity modification

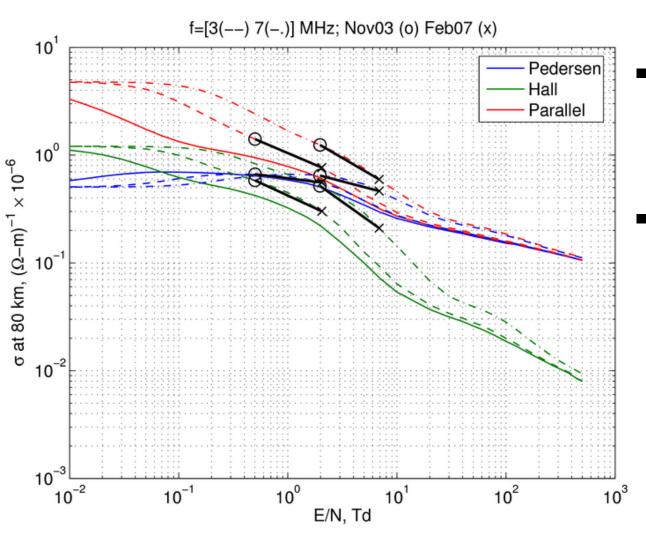




- Pedersen conductivity is increased
- Parallel conductivity is decreased



Conductivity as a function of E/N (x-mode, h=80 km, f=0,3,7 MHz)

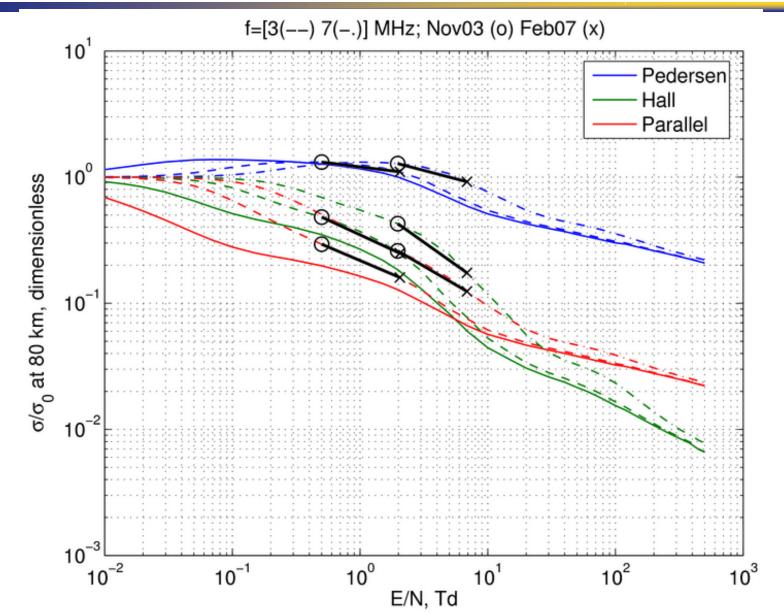


- Solid line shows conductivity modifications by DC field
- Black intervals
 connect the
 conductivities
 modified by
 maximum
 HAARP heating
 before and after
 upgrade



Relative change of conductivity $\sigma(E)/\sigma(E=0)$





20





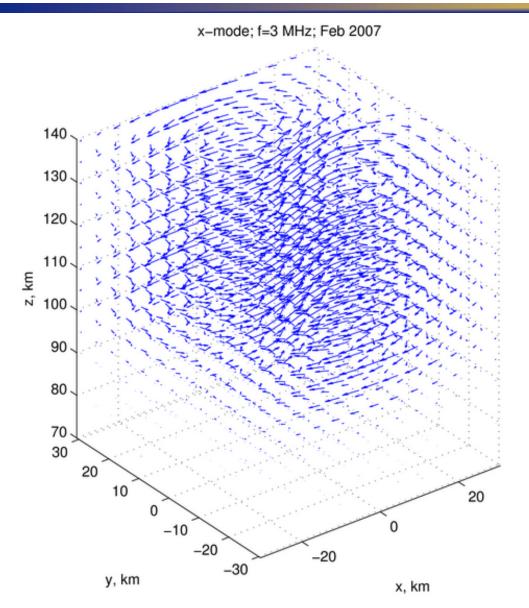
- In most previous works, it is assumed that the electrojet field E_{ej}=const => inaccurate at low frequencies (no account for the accumulation of charge)
- We assume static current, i.e.

$$\vec{J} = -\vec{\sigma}\nabla \varphi$$
$$\nabla \cdot \vec{J} = 0$$



3D stationary ΔJ



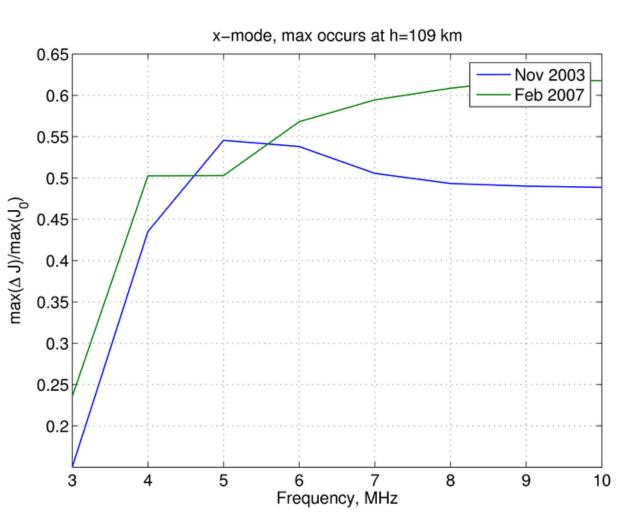


- Vertical B
- Ambient E is along x
- Ambient current is mostly along y
- Models with ∆E=0 do not consider closing side currents
- max ∆J/J₀~0.3 for this case



Calculated $\Delta J/J_0$ for various frequencies





- Range 70-130 km
- Modified region radius ~10 km before upgrade and ~5 km after upgrade
- Calculated maximum current and its modification occur at ~109km





- Our model includes both:
 - Non-Maxwellian electron distribution
 - Self-absorption
- Maxwellian electron distribution models, which calculate ΔT_e , cannot account for the nonlinear T_e saturation.
- The non-Maxwellian model allows to calculate processes for which high-energy tail of the electron distribution is important, such as:
 - optical emissions
 - breakdown processes.





- Electrojet current modulation in nonstatic case
- ELF/VLF emission
- ELF/VLF wave propagation along the geomagnetic field line